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The Crystal Structure of Sodium Trihydrogen Selenite

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Sodium trihydrogen selenite is ferroelectric below -79 °C. At room temperature this compound crystallizes in the monoclinic space group $P2_1/n$ with two formula units in a unit cell of dimensions a=10.35, b=4.84, c=5.84 Å and $\beta=89^{\circ}$. The structure was solved by the interpretation of an unsharpened three-dimensional Patterson synthesis and was refined from photographic X-ray data by the structure-factor least-squares method to an R index of 0.134 for 491 independent observed reflexions. In the structure, the selenite ions have idealized $3m(C_{3v})$ symmetry. The coordination of oxygen atoms around the sodium ion is a slightly distorted octahedron. The selenite ions are held together by a system of strong hydrogen bonds. These hydrogen bonds form closed loops and may work as switching mechanisms in the ferroelectric phase.

Introduction

Since the discovery of ferroelectricity in sodium trihydrogen selenite, $NaH_3(SeO_3)_2$, below $-79^{\circ}C$ by Pepinsky & Vedam (1959), this compound has been subjected to various physical investigations by many workers. Blinc and his co-workers carried out nuclearmagnetic and electron-paramagnetic resonance studies on NaH₃(SeO₃)₂ and the deuterated salt NaD₃(SeO₃)₂ (Blinc & Pintar, 1961; Blinc, Poberaj, Sehara & Stepinsnik, 1966). The dielectric behaviour of the two compounds as a function of temperature has been reported (Blinc, Jovanovic, Levstik & Prelesnik, 1965; Gavrilova-Podol'skaya, Yudin & Lundin, 1965). The Raman and infrared spectra of sodium trihydrogen selenite have also been analysed (Bazhulin, Mysasnikova & Rakov, 1964; Khanna, Decius & Lippincott, 1965). However, no detailed structural information on it is so far available. Hence it was thought worth while to analyse its structure to understand the structural characteristics that are responsible for the occurrence of ferroelectricity in this crystal. The determination of the structure at room temperature is reported here. Preliminary crystallographic data on this compound have already been published (Vijayan, 1967).

The corresponding lithium salt, $LiH_3(SeO_3)_2$, is a room-temperature ferroelectric and its structure has

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been solved (Vedam, Okaya & Pepinsky, 1960). However, it is not isomorphous with sodium trihydrogen selenite.

Experimental

Transparent plate-shaped crystals of sodium trihydrogen selenite, elongated along the *b* axis, were grown by slow evaporation at room temperature of a saturated aqueous solution of selenium dioxide and sodium carbonate in stoichiometric molar ratio. The space group and unit-cell dimensions were determined from oscillation and Weissenberg photographs taken about crystallographic axes with Cu $K\alpha$ radiation. The density of the crystal was measured by flotation in a liquid mixture of methylene iodide and bromoform.

Crystal data

NaH₃(SeO₃)₂ F.W. 279.93 Monoclinic $P2_1/n$ $a=10.35\pm0.02$, $b=4.84\pm0.02$, $c=5.84\pm0.02$ Å, $\beta=89\pm1^{\circ}$. $D_m=3.20\pm0.01$ g.cm⁻³, U=292.37 Å³, Z=2, $D_x=3.18$ g.cm⁻³. Absorption coefficient for X-rays, ($\lambda=1.5418$) $\mu=185.5$ cm⁻¹.

Intensity data were collected for reciprocal levels hkl, k = 0-4, from a nearly cylindrical specimen of mean radius 0.016 cm using nickel-filtered copper radiation. hk0 zonal data were also recorded with a specimen of mean radius 0.015 cm. The well known multiple-film equi-inclination Weissenberg method was used for data

Table 1. Fractional positional coordinates and isotropic temperature factors

The standard deviations ($\times 10^5$) of the positional coordinates are given in parentheses.

	x	У	Z	В
Se	0.16907 (29)	0.11496 (70)	-0.00682(53)	1·023 Ų
Na	0.50000 (0)	0.00000 (0)	0.00000 (0)	2.595
0(1)	0.31108(208)	0.30205 (475)	0.06351 (361)	1.534
$\tilde{O}(2)$	0.05123(224)	0.30523 (518)	0.12641 (405)	2.028
Õ(3)	0.14103 (231)	0.20148 (532)	-0·29748 (417)	2.139

collection and the intensities were estimated visually. Out of the 715 independent reflexions in the copper sphere, 638 were recorded, of which 491 were in the measurable range. The data were corrected for Lorentz

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and polarization factors. Absorption corrections were applied assuming the specimens to be perfect cylinders (*International Tables for X-ray Crystallography*, 1959). However, it may be mentioned that the specimens were

Table 2. Observed and calculated structure factors

$ \begin{array}{c} \textbf{k} \\ \textbf{F}_{0} \\ \textbf{c}_{24} \\ \textbf{c}_{25} \\ \textbf{c}_{24} \\ \textbf{c}_{24} \\ \textbf{c}_{25} \\ \textbf{c}_{24} \\ \textbf{c}_{24} \\ \textbf{c}_{25} \\ \textbf{c}$			K For
$\begin{array}{c} 2 & 4 \\ 70 \\ 70 \\ 70 \\ 70 \\ 70 \\ 70 \\ 70 \\ $	LUALDA A A A A A A A A A A A A A A A A A A	140071890000770077 553346573605789445514441336 4202070	$\begin{array}{c} 1 & -1055 \\ -1046 \\ 2 & 7711 \\ 2 & 7711 \\ 6 & 633 \\ 4 & 5673 \\ 5 & 7772 \\ 5 & 6873 \\ 5 & 7772 \\ 5 & 6873 \\ 5 & 7772 \\ 5 & 7772 \\ 5 & 7772 \\ 5 & 7772 \\ 5 & 7772 \\ 5 & 7772 \\ 5 & 7772 \\ 7 & 777$

not perfectly cylindrical and this is a limiting factor for the accuracy of thermal parameters derived from this analysis.

Structure analysis

The multiplicity of general positions in the space group is fourfold, and as there are only two formula units in the elementary cell, the sodium ion and one of the hydrogen atoms should occupy twofold special positions at the inversion centres. The rest of the atoms are at general positions.

The structure was solved by the direct interpretation of a three-dimensional unsharpened Patterson function. The atomic coordinates thus derived were confirmed by a selenium and sodium phased three-dimensional Fourier synthesis and projection minimum maps. The positional and individual isotropic temperature parameters were refined in several cycles of structurefactor least-squares calculations to an R index, defined as

$$\Sigma ||F_o| - |F_c|| / \Sigma |F_o| ,$$



Fig. 1. Superposed sections of the final electron density distribution through atomic centres parallel to (010). The contours start at 2 e.Å⁻³ and are at intervals 2 e.Å⁻³ for oxygen, 3 e.Å⁻³ for sodium and 10 e.Å⁻³ for selenium atoms.



Fig. 2. Dimensions of the selenite ion. Standard deviations are given in parentheses.

of 0.134 for all the observed reflexions. These refinement cycles were done on the Elliott 803-B Computer at the Hindustan Aeronautics Ltd., Bangalore, using a block-diagonal SFLS program devised by Dr G.A. Mair of the Royal Institution, London. The quantity minimized in the program is $\Sigma w|kF_o - F_c|^2$. In the final cycle, all the indicated shifts were less than half the corresponding standard deviations, the maximum and the average shifts being 0.38 σ and 0.13 σ respectively. The weighting scheme employed in these calculations was of the form

$$\psi w = 1 / \left\{ 1 + \left(\frac{kF_o - b}{a} \right)^2 \right\}^{1/2}$$

with a = 32, b = 20 for $k = \frac{1}{2}$. The atomic scattering factors used were those given by Cromer & Waber (1965).

The final positional and isotropic thermal parameters of the atoms in the asymmetric part of the unit cell are given in Table 1. Table 2 lists the observed and calculated structure factors (unobservables excluded). Superposed sections through atomic centres of the final electron density distribution are shown in Fig. 1. A final $F_o - F_c$ synthesis was computed, but the positions of the hydrogen atoms could not be located.

Discussion

Dimensions of the SeO_3^{2-} ion

The dimensions of the selenite ion in the structure are shown in Fig. 2.

The pyramidal selenite ion has an idealized $3m(C_{3v})$ symmetry. However, the departure from this symmetry is statistically significant. The Raman spectrum of this compound also shows that, in the crystalline state, the symmetry of SeO_3^{2-} ion is decreased from 3m (Bazhulin *et al.*, 1964). Two Se-O bonds in the ion are almost equal in length whereas the length of the third is significantly different ($\Delta > 2\sigma$) from the other two. The O-Se-O angles are unequal though comparable in magnitude. The selenium atom is 0.77 Å away from the plane of the three oxygen atoms. The normal to this plane makes an angle of 12° with the crystallographic *b* axis.

It may be mentioned that Vedam *et al.* (1960) have quoted the dimensions of the selenite ion in sodium trihydrogen selenite from a private communication. But neither the details nor the other results of this work have been published so far. However, the values quoted compare well with those obtained from the present analysis.

Coordination around the Na⁺ ion

The sodium ion in the structure occupies a centre of symmetry at $\frac{1}{2}$, 0, 0. The coordination polyhedron of selenite oxygen atoms around this ion is a slightly distorted octahedron. The dimensions of this octahedron are listed in Table 3 and the environment of the sodium ion, as seen down the *b* axis, is shown in Fig. 3.

 Table 3. Dimensions of the coordination

 octahedron around the Na⁺ ion

A x, y, z		
$B = \frac{1}{2} - x, (\frac{1}{2} + y) - \frac{1}{2}$	$1, \frac{1}{2} - z$	
$C \frac{1}{2} + x, \frac{1}{2} - y, \frac{1}{2} +$	z	
D - x + 1, -y, -	z	
$E = \frac{1}{2} + x, = \frac{1}{2} - y, = \frac{1}{2} - \frac{1}{2}$	(+z) - 1	
$F = \frac{1}{2} - x, (\frac{1}{2} + y) - \frac{1}{2}$	$1, (\frac{1}{2} - z) - z$	- 1
Na - O(1)A	2 ·464 A	A (×2)
Na - O(2)B	2.427	(×2)
Na - O(3)C	2.383	$(\times 2)$
O(1)A - O(3)C	3.524	$(\times 2)$
O(3)C-O(2)E	3.542	(×2)
O(2)E-O(1)A	3.570	$(\times 2)$
O(1)A - O(2)B	3.344	$(\tilde{x} \tilde{2})$
O(2)B-O(3)C	3.254	$(\times 2)$
O(3)C = O(1)D	3.329	$(\times 2)$
O(1)A - Na - O(2)B	86·3°	$(\times 2)$
O(2)B-Na-O(3)C	85.2	$(\times 2)$
O(3)C - Na - O(1)D	86.7	$(\times 2)$
O(2)B-Na-O(1)D	93.8	$(\times 2)$
O(1)D - Na - O(3)F	93.3	$(\times 2)$
O(3)E - Na - O(2)B	94.9	(\hat{x}_{2})

Crystal structure and hydrogen bonding

The contents of the unit cell as viewed in the [010] direction are shown in Fig.4. Fig.5 shows a view of the structure projected on the (001) plane.

The selenite ions in the structure are held together by a three-dimensional network of strong hydrogen bonds which are indicated by broken lines in Fig.5. All the oxygen atoms in the structure take part in hydrogen bond formation either as donors or as acceptors. The two crystallographically non-equivalent hydrogen-bonded $O \cdots O$ distances and the relevant angles are given in Table 4. These bonds may be considered strong judging by the short $O \cdots O$ distances. It may be mentioned that hydrogen bonds of comparable lengths are found in other hydrogen-bonded ferroelectrics like potassium dihydrogen phosphate and lithium trihydrogen selenite (Frazer & Pepinsky, 1953; Bacon & Pease, 1953, 1955; Vedam et al., 1960). The hydrogen bonds involving O(1)A, O(3)B and O(2)A, O(2)C are inclined to the crystallographic b axis at 40° and 42° respectively. The oxygen atoms O(2)A and O(2)C are related by a centre of inversion and hence the bond connecting them should be symmetrical. The length of this bond is, however, greater than that normally found for symmetric hydrogen bonds.

Table 4. Hydrogen bond distances and relevant angles

$\begin{array}{c} A \ x, y, z \\ B \ \frac{1}{2} - x, \ \frac{1}{2} + y, \ (\frac{1}{2} - C \ -x, \ -y + 1, \ -y \end{array}$	(z) - 1
$\begin{array}{c} O(1)A\cdots O(3)B\\ O(2)A\cdots O(2)C \end{array}$	2·524 Å 2·629
$SeA-O(1)A\cdots O(3)B$ $SeB-O(3)B\cdots O(1)A$ $SeA-O(2)A\cdots O(2)C$ $SeC-O(2)C\cdots O(2)A$	113·7° 112·0 115·0

Khanna *et al.* (1965) have explained the infrared spectrum of sodium trihydrogen selenite in the paraelectric phase assuming an even proton-density distribution along the hydrogen bond length. This might be due, as observed in potassium dihydrogen phosphate (Bacon & Pease, 1953, 1955), to a symmetric double minimum potential or a broad single minimum potential for the proton. The former situation may be described physically as resulting from a statistical disorder of the proton in the structure whereas the latter as a result of a high thermal-vibration amplitude for the proton. This probably explains the comparatively large length of the symmetrical hydrogen bond between O(2)A and O(2)C.

Dielectric studies show that the ferroelectric transition point is altered from -79 °C for NaH₃(SeO₃)₂ to -29 °C for NaD₃(SeO₃)₂ (Blinc *et al.*, 1965; Gavrilova-Podol'skaya *et al.*, 1965). This result indicates that hydrogen bonds play a finite role in the ferroelectric



Fig. 3. Environment of the sodium ion as seen down the b axis.



Fig.4. The contents of the unit cell as viewed in the [010] direction.



Fig. 5. A view of the structure projected onto the (001) plane. Hydrogen bonds are indicated by broken lines. The code of symmetry-related positions is as follows:

 $\begin{array}{ll} & x, y, z \\ \text{II} & \frac{1}{2} - x, \frac{1}{2} + y, (\frac{1}{2} - z) - 1 \\ \text{III} & \frac{1}{2} + x, (\frac{1}{2} - y) + 1, (\frac{1}{2} + z) - 1 \\ \text{IV} & -x + 1, -y + 1, -z - 1 \\ \text{V} & \frac{1}{2} + x, \frac{1}{2} - y, (\frac{1}{2} + z) - 1 \\ \text{VI} & \frac{1}{2} - x, (\frac{1}{2} - y) - 1, (\frac{1}{2} - z) - 1. \end{array}$

transition in this compound. In the ferroelectric phase the centre of symmetry vanishes and consequently, the selenite ions related by it become non-equivalent. The infrared studies quoted above show that the transition to the ferroelectric phase is accompanied by an ordering of the protons in the structure. The disposition of hydrogen bonds in the structure is such that they form closed loops (Fig. 5). All these results suggest that the system of hydrogen bonds works, as in the case of potassium dihydrogen phosphate, as a switching mechanism in the ferroelectric phase, *i.e.* the roles of the donor and acceptor oxygen atoms are reversed when the polarity of the applied electric field is reversed.

All the oxygen atoms in the structure are coordinated to one or another of the sodium ions. The coordination octahedron possesses a centre of symmetry with the sodium ion at the centre. In the ferroelectric phase, the centre of symmetry vanishes and the sodium ion may occupy an asymmetric site giving rise to a net dipole moment. Alternatively, in the ferroelectric phase the oxygen octahedron may undergo reversible distortions in an applied electric field. It is probable that these distortions contribute substantially to the ferroelectric behaviour of the crystal. However, definite conclusions can be drawn only after the crystal structure is analysed in the ferroelectric phase also.

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